Circular Polarized Carbon-NanotubePatch Antenna embedded in Superstrates Anisotropic Media

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Abstract: The current distribution on surface of the Patch Antenna is radiating in the presence of the Superstrates layered anisotropic medium. The electric and magnetic field in the upper half-space are formulated using Dyadic Green Function in terms of Sommerfeld-Weyl-type integrals. The stationary phase method is used to obtain far-field expressions. Carbon-Nanotube Composite (CNT) as the radiating element for Star Patch Antenna is used. .Return loss, Axial Ratio (<3 dB Circular Polarized), Gain and Directivity parameters of Carbon-Nanotube Star Patch Antenna are compared for different anisotropic materials.

Keywords: Anisotropic, patch antenna, Moment Method, Dyadic Green Function. Carbon-Nanotube Composite

I. Introduction

Analyzing the antenna parameters with anisotropic media was appeared in many works [1-4]. The affect of anisotropic ratio on the far-field radiation pattern was investigated ([5]. Because of cost, weight, fabrication procedure it limited the use of copper or metals for patch or ground plane. Several composite materials as replacement for metals were studied in many recent studies [6-8]. Due to electrical properties and favorable mechanical Carbon Nanotube Composites have been used for electromagnetic applications [9-12]. In this study, using the dyadic Green's function formulation the electric and magnetic fields in the upper half-space due to the current distribution on the Star Patch are obtained [13]. The Sommerfeld-Weyl-type integrals of electric and magnetic fields are calculated in the radiation zone or far-field region [13].





II. Expressing The Electric and Magnetic Fields 2.1 Electric and Magnetic Fields in the Upper-Half-space

In the upper-half space (region 0), the far-field satisfies the following relations:

$$E_{0\theta} = -\frac{E_{0z}}{\sin\theta} \tag{2.1}$$

$$H_{0\phi} = \frac{1}{\eta_0} E_{0\theta} \tag{2.2}$$

for TM waves, and

$$H_{0\theta} = -\frac{H_{0z}}{\sin\theta}$$
(2.3)

$$E_{0\phi} = -\eta_0 H_{0\theta} \tag{2.4}$$

for TE waves, and

$$\eta_0 = \sqrt{\frac{\mu_0}{\varepsilon_0}} \tag{2.5}$$

The electric field E_{0z} can be obtained as

$$E_{0z} = \hat{z} \cdot \overline{E}_0 = \hat{z} \cdot i\omega \iint_{S'} dr' \,\overline{\overline{G}}_{01}(\bar{r}, \bar{r}') \cdot \,\overline{J}_s(\bar{r}\,') = i\omega \iint_{S'} dx' \, dy' \, G_{01}^{(z,x)}(\bar{r}, \bar{r}\,') \, J_x(x', y')$$
(2.6)

where

$$G_{01}^{(z,x)}(\bar{r},\bar{r}') = \hat{z} \cdot \overline{\overline{G}}_{01}(\bar{r},\bar{r}') \cdot \hat{x}$$

$$(2.7)$$

 $\overline{\overline{G}}_{01}(\overline{r},\overline{r}')$ can be obtained as shown in [13] (See Fig. 1). For anisotropic medium:

$$G_{01}^{(z,x)}(\bar{r},\bar{r}') = \int_{-\infty}^{\infty} dk_x \, dk_y \, e^{i\bar{k}_s \cdot (\bar{r}_s - \bar{r}'_s)} \, g_{01}^{(z,x)}(\bar{k}_s, z, z')$$
(2.8)

where

$$g_{01}^{(z,x)}(\bar{k}_s,z) = \frac{-i\mu_0}{8\pi^2} \frac{r_1^{TM}}{k_0} X_{\cup 0,1}^{TM} e^{ik_{0z}z} e^{ik_{1z}^{(e)}h_1} \frac{k_x}{k_1} \left(1 - R_{\cap 1}^{TM} e^{ik_{1z}^{(e)}h_1}\right)$$
(2.9)

where $\,z^{\,\prime}\,$ is substituted by ($-h_{\!_1}$).

Thus, the z-component of the electric field in region (0) is given by

$$E_{0z} = (i) \frac{\omega \mu_0}{4\pi} \cdot \iint_{S'} dx' \, dy' \, J_x(x', y') \cos \phi \cdot \int_0^\infty dk_\rho \, k_\rho^2 g_E(k_\rho) \, e^{ik_{0z}z} J_1(k_\rho \, |\overline{\rho} - \overline{\rho}'|) (2.10)$$

as Sommerfeld-type integral [12].

The magnetic field in the upper half-space can be obtained from Maxwell's equations

$$\overline{H}_0 = \frac{1}{i\omega\mu_0} \nabla \times \overline{E}_0 \tag{2.11}$$

That is

$$\overline{H}_{0}\left(\overline{r}\right) = \frac{1}{\mu_{0}} \iint_{S'} dx' \, dy' \, \nabla \times \overline{\overline{G}}_{01}(\overline{r}, \overline{r}') \cdot \, \overline{J}_{s}(x', y') \tag{2.12}$$

where \equiv

$$\mathbf{G}_{01}(\bar{r},\bar{r}') = \frac{i\mu_{0}}{8\pi^{2}} \int_{-\infty}^{\infty} d\bar{k}_{s} e^{i\bar{k}_{s}\cdot(\bar{r}_{s}-\bar{r}_{s}')} e^{ik_{0z}z} \left\{ \frac{r_{1}^{TE}}{k_{1z}^{(h)}} X_{\cup0,1}^{TE} e^{ik_{1z}^{(h)}h_{1}} \hat{h}(k_{0z}) \left[\hat{h}(k_{1z}^{(h)}) + R_{\cap 1}^{TE} e^{ik_{1z}^{(h)}h_{1}} \hat{h}(-k_{1z}^{(h)}) \right] \right\} + \frac{r_{1}^{TM}}{k_{1z}^{(e)}} X_{\cup0,1}^{TM} e^{ik_{1z}^{(e)}h_{1}} \hat{v}(-k_{1z}^{(h)}) \right] \qquad (2.13)$$

Hence, the magnetic field H_{0z} can be obtained as

$$H_{0z} = \frac{-1}{4\pi} \iint_{S'} dx' \, dy' \, J_x(x', y') \sin \phi \cdot \int_0^\infty dk_\rho \, k_\rho^2 g_H(k_\rho) \, e^{ik_{0z}z} \cdot J_1\left(k_\rho \left|\overline{\rho} - \overline{\rho}'\right|\right) \tag{2.14}$$

as Sommerfeld type integral.

2.2 Far-Field Expressions

The z- component of the electric field in the upper half-space given by Eq. 2.10, can be written as

$$E_{0z}(\bar{r}) = \frac{\iota \omega \mu_0}{4\pi} \iint_{S'} dx' \, dy' \, J_x(x', y') \cos \phi \, I_E$$
(2.15)

where

$$I_E = \int_0^\infty dk_\rho \ k_\rho^2 \ g_E(k_\rho) \ e^{ik_{0z}z} \ J_1(k_\rho |\overline{\rho} - \overline{\rho}'|)$$
(2.16)

which can be put in the form

$$I_E = \frac{1}{2} \int_{-\infty}^{\infty} dk_{\rho} k_{\rho}^2 g_E(k_{\rho}) e^{ik_{0z}z} H_1^{(1)}(k_{\rho}|\overline{\rho}-\overline{\rho}'|)$$

$$(2.17)$$

Using the large argument approximation for the Hankel function, the rapidly varying part of the integrand in Eq. 2.17 is asymptotic to the phase factor $e^{ik_{\rho}|\bar{\rho}-\bar{\rho}'|+ik_{0z}z}$. Thus, the location of the stationary phase point is given by

$$\frac{\partial}{\partial k_{\rho}} \left. i \left(k_{\rho} \left| \overline{\rho} - \overline{\rho}' \right| + k_{0z} z \right) \right|_{k_{\rho 0}} = 0$$
(2.18)

The solution is given by $k_{\rho_0} = k_0 \sin\theta, \qquad k_{0z0} = k_0 \cos\theta$ (2.19)

with
$$\theta = \sin^{-1} \frac{|\overline{\rho} - \overline{\rho}'|}{|\overline{r} - \overline{\rho}'|}$$
, (see Fig. 2.2).



Fig. 2, Relation between ρ and θ

This leads to the following expression for I_E

$$I_E \sim -i g_E(k_{\rho_0}) k_{\rho_0} e^{-i\pi/2} k_{0z0} \frac{e^{ik_0|\bar{r}-\bar{\rho}'|}}{|\bar{r}-\bar{\rho}'|} , r \to \infty$$

$$(2.20)$$

But

$$\left| \overline{r} - \overline{\rho}' \right| \cong r - \hat{r} \cdot \overline{\rho}',$$

thus

$$I_E \sim k_0^2 \sin\theta \cos\theta g_E(k_{\rho_0}) \ e^{-i\bar{k_0}\cdot\bar{\rho}'} \ \frac{e^{i\kappa_0 r}}{r} \ , \ r \to \infty$$
(2.21)

Thus,

$$E_{0z}(\bar{r}) \sim \frac{-i\omega\mu_0}{4\pi} k_0^2 \sin\theta \,\cos\theta \,\cos\varphi_{BE}(k_{\rho_0}) \,\frac{e^{ik_0r}}{r} \times \iint dx' \,dy' \,J_x(x',y') \,e^{-i\bar{k_0}\cdot\bar{\rho}'} \tag{2.22}$$

Following similar procedure, we can get for the magnetic field in the upper half-space as $r \rightarrow \infty$, the following expression

$$H_{0z}(\bar{r}) \sim \frac{i}{4\pi} k_0^2 \sin\theta \cos\theta \sin\Phi g_H(k_{\rho_0}) \frac{e^{ik_0r}}{r} \times \iint dx' \, dy' \, J_x(x',y') \, e^{-i\bar{k_0}\cdot\bar{\rho}'} \tag{2.23}$$

Thus, using Eq. 2.1 - Eq. 2.4, the field patterns are readily obtained

III. Results And Discussion

CNT Star patch numerical results embedded in anisotropic superstrate structures are presented. Figure 2, shows the geometry of CNT Star Patch Antenna. The effects of anisotropy of Superstrate layers in S-parameters, Axial Ratio, Directivity and gain are investigated. To simulate the Far Field Radiation of CNT Star Patch, a Matlab program is written. For Validation propose simulation software CST STUDIO SUITE® is used. The compered result between the Two programs was good similarity. A microstrip CNT Star patch of length Lp = 42.12 mm and width W = 42.12mm on the top of RTDURIOD 5870 an isotropic substrate of permittivity $\varepsilon_{\rm T} = 2.33$, and thickness hs = 1.5mm. The Ground plate of square shape with length L= 52mm. The CTN Star patch is fed with coaxial probe. The following anisotropic substrates are used: pirolytic boron nitride, or pbn $\varepsilon_{xx} = 5.12; \varepsilon_{zz} = 3.4$, sapphire ($\varepsilon_{xx} = 9.4; \varepsilon_{zz} = 11.6$), Epsilam10 ($\varepsilon_{xx} = 13; \varepsilon_{zz} = 10.6$). For $\varepsilon_r = 1$ (vacuum or Air) use Foam flakes $\varepsilon_r = 1.1$, which is good approximation.

3.1 Input Impedance of Star patch with superstrate anisotropic layer

The Anisotropic Superstrate effect for CNT Star Patch Antenna is conducted. The Superstrate Layer assumed to be anisotropic media with permittivity pirolytic boron nitride (pbn) $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$. The E-Plane Radiation of CNT Star Patch antenna fed by Coaxial Cable was shown in Fig. 3.



Fig. 3, E – Plane radiation of CNT Star Patch Antenna

The CST Simulating result almost identical to the theoretical model. **3.2 Farfield Directivities and gainsof CNT Star patch without any Superstrate** A CNT Star Patch without any Superstrate is considered. Figure 4,show, the circuit of CNT Star patch.



Fig. 4, CNT Star Patch Antenna

Table 1, Directivities and Gains of CNT Star Patch Antenna without Superstrate

F	Directivity	Gain
GHz	dBi	dBi
5.5	3.95	2.11
8.5	5.94	2.48
10.5	3.65	0.442
12.5	1.39	-0.512
14.5	4.49	1.49
17.5	4.54	2.3
19.5	4.96	3.02
21.5	3.19	1.69
23.5	4.51	2.71
25.5	9.53	8.06



Farfield Directivity Abs (Phi=90)







Design Criteria's for CNT Star Patch Antenna are :

- 1) $S11 \le -10 \, dB$
- 2) Axial Ratio $\leq 3 \text{ dB}$
- 3) Gain should be maximum
- 4) Directivity should be maximum

Frequency GHz	S11 dB	Axial Ratio dB	Gain dB	Directivity dB
7.517	-22.75	6.3	4.4	5.53
12.123	-16.6	8.02	3.79	4.43
14.279	-18.47	0.661	4.52	5.75
16.582	-27.19	1.456	8.66	9.26
18.15	-11.52	3.49	6.32	7.29
20.894	-24.04	1.099	3.7	3.94
24.961	-13.32	2.704	8.68	8.95

 Table 2, S11, Axial Ratio, Gain and Directivities for best Design of CNT Star Patch Antenna without Superstrate



Fig. 8, Axial Ratio, Gain and Directivities for best Design of CNT Star Patch Antenna without Superstrate

CST Studio Suite designs the CNT Sstar patch (Fig. 4). Fig. 7. Shows S11 for this configuration. It is clear from figure 5, star patch directivity is 11.2 dBi. The gain of that patch is 9.69 dBi as shown in Fig. 6. From Table 2 and Figure 8, the best Design parameters are : F = 16.582 GHzS11 = -27.19 dB

Axial Ratio = 1.456 dBGain = 8.65 dBDirectivity = 9.26 dB

3.3 Farfield Directivities and gains f CNT Star patch with one Superstrate layer



Fig. 9. CNT Star Patch Antenna with one anisotropic superstrate Layer, $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$

F	Directivity	Gain
GHz	dBi	dBi
1.5	4.54	-2.46
2.5	7.09	-1.36
4.5	4.3	2.23
7.5	4.65	2.84
11.5	12.2	11
17.5	5.94	4.86
19.5	5.85	4.94
21.5	6.85	5.44

Table 3 Directivities and gains of CNT Star Patch Antenna with one Anisotropic Superstrate layer , $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$



Fig. 10. Directivity of CNT Star Patch at F=11.5 GHz with one Anisotropic Superstrate layer, , $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$



Fig 11. Gain of CNT Star Patch at F=11.5 GHz with one Anisotropic Superstrate layer, , $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$



Fig. 12. S11 of CNT Star Patch Antenna with one Anisotropic Superstrate layer, , $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$

T insolution is superstruct rayer, $t_{22} = 5.12$, $t_{22} = 5.14$.				
Frequency GHz	S11 dB	Axial Ratio dB	Gain dB	Directivity dB
5.399	-10.89	6.03	3.47	4.99
9.116	-20.66	3.901	4.38	5.62
10.691	-12.63	4.284	-2.63	-1.11
14.849	-13.63	1.574	4.48	5.60
24.509	-10.68	4.186	7.88	8.89

Table 4, S11, Axial Ratio, Gain and Directivities for best Design of CNT Star Patch Antenna with one Anisotropic Superstrate layer, $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$.



Fig. 13, Axial Ratio, Gain and Directivities for best Design of CNT Star Patch Antenna with one Anisotropic Superstrate layer, $\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$.

Fig. 9. Shows the Directivity for CNT Star Patch with pirolytic boron nitride, or pbn($\varepsilon_{xx} = 5.12$; $\varepsilon_{zz} = 3.4$.) Superstrate is greater than CNT Star patch without Superstrate. In addition, the Gain also increased if adding a Anisotropic Superstrate. Fig. 12, shows S11.

From Table 4 and Fig. 13, the best Design parameters are : F = 14.849 GHz S11 = -13.63 dBAxial Ratio = 1.574 dB Gain = 4.48 dB Directivity = 5.66 dB

3.4 Farfield Directivities and gainsof CNT Star patch with two Superstrate layers



Fig. 14. CNT Star Patch Antenna with two anisotropic Superstrate Layers, $\varepsilon \mathbf{1}_{xx} = 9.4$; $\varepsilon \mathbf{1}_{zz} = 11.6$ (layer 1) and $\varepsilon \mathbf{2}_{xx} = 13$; $\varepsilon \mathbf{2}_{zz} = 10.6$ (layer 2)

Table 5. Directivities and gains of CNT Star Patch Antenna with two Anisotropic Superstrate Layers, $\varepsilon 1_{xx}$	=
9.4; $\varepsilon 1_{zz} = 11.6$ (layer 1) and $\varepsilon 2_{xx} = 13$; $\varepsilon 2_{zz} = 10.6$ (layer 2).	

F	Directivity	Gain
GHz	dBi	dBi
4.5	3.93	3.37
5.5	4.31	3.11
7.5	4.39	3.61
9.5	4.65	3.87
11.5	3.97	2.95
13.5	3,97	2.64
15.5	7.72	4.78
17.5	4.98	3.29
19.5	13.2	8.56
19.75	14.8	8.2





Fig. 15. Directivity of CNT Star Patch at F=19.75 GHz with two Anisotropic Superstrate Layers, $\varepsilon \mathbf{1}_{xx} = 9.4$; $\varepsilon \mathbf{1}_{zz} = 11.6$ (layer 1) and $\varepsilon \mathbf{2}_{xx} = 13$; $\varepsilon \mathbf{2}_{zz} = 10.6$ (layer 2).



Fig. 16. Gain of CNT Star Patch at F=19.75 GHz with two Anisotropic Superstrate Layers, $\varepsilon 1_{xx} = 9.4$; $\varepsilon 1_{zz} = 11.6$ (layer 1) and $\varepsilon 2_{xx} = 13$; $\varepsilon 2_{zz} = 10.6$ (layer 2).



Fig. 17. S11 of CNT Star Patch Antenna with two Anisotropic Superstrate Layers, $\varepsilon \mathbf{1}_{xx} = 9.4$; $\varepsilon \mathbf{1}_{zz} = 11.6$ (layer 1) and $\varepsilon \mathbf{2}_{xx} = 13$; $\varepsilon \mathbf{2}_{zz} = 10.6$ (layer 2).

Frequency	S11	Axial Ratio	Gain	Directivity
GHz	dB	dB	dB	dB
4.289	-19.39	19.466	3.27	4.56
5.088	-20.93	1.105	2.98	4.36
5.836	-15.697	1.092	3.98	5.42
7.672	-12.699	1.05	1.15	1.49
7.808	-11.45	1.654	2.21	2.99
7.995	-13.29	1.246	2.67	3.62
8.675	-23.14	1.236	0.722	1.61
10.273	-11.4	1.43	3.31	4.18
11.701	-15.5	1.374	10.3	14
12.364	-12.57	1.46	6.07	8.88
12.84	-12.08	1.118	4.48	6.6
13.027	-10.84	1.1657	1.72	2.69
13.673	-11.97	2.06	3.73	5.99
14.251	-13.19	1.27	4.15	6.16
14.438	-12.92	1.11	5.29	7.48
15.628	-21.999	1.058	5.64	8.27
16.478	-12.07	1.59	2.68	3.92
17.107	-10.46	1.236	2.75	3.8
17.6	-13.92	1.244	1.85	2.79
18.224	-12.106	1.44	1.04	1.63
18.416	-14.59	1.125	3.41	5.29
19.028	-12.106	1.135	2.33	3.42
19.436	-10.35	1.57	3.94	6.5
19.776	-11.6	1.22	1.03	1.85
20.949	-1.994	1.22	7.12	11.5

Table 6, S11, Axial Ratio, Gain and Directivities for best Design of CNT Star Patch Antenna with two Anisotropic Superstrate Layers, $\varepsilon 1_{xx} = 9.4$; $\varepsilon 1_{zz} = 11.6$ (layer 1) and $\varepsilon 2_{xx} = 13$; $\varepsilon 2_{zz} = 10.6$ (layer 2).



Fig. 18, Axial Ratio, Gain and Directivities for best Design of CNT Star Patch Antenna with two Anisotropic Superstrate Layers, $\varepsilon 1_{zz} = 9.4$; $\varepsilon 1_{zz} = 11.6$ (layer 1) and $\varepsilon 2_{zz} = 13$; $\varepsilon 2_{zz} = 10.6$ (layer 2).)

In this configuration, sapphire ($\varepsilon_{xx} = 9.4$; $\varepsilon_{zz} = 11.6$), for Superstrate 1 is used. While), Epsilam-10 ($\varepsilon_{xx} = 13$; $\varepsilon_{zz} = 10.6$).is used for superstrate 2. From Fig. 15. Directivity was increased to 14.8 dBi, and the gain was 8.2 at frequency fr = 19.75 GHz From Table 6 and Fig. 18, the best Design parameters are : F = 11.701 GHz S11 = -15.5 dB Axial Ratio = 1.374 dB Gain = 10.3 dB Directivity = 14 dB

IV. Conclusion

This paper shows a high gain and directivity for CNT Star Patch Antenna embedded in multilayers anisotropic media. Also, it shows, how the anisotropic phenomena can improve the performance of the patch antenna.

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